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# Improvement in the thermal shock resistance of alumina through the addition of submicron-sized aluminium nitride particles

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## Abstract

The possibility of improving the thermal shock behaviour of alumina under mild thermal conditions due to the addition of submicron-sized and homogeneously dispersed AlN particles is studied. Processing conditions were adjusted in order to avoid the formation of AlON and, consequently, a decrease in the thermal conductivity of the composite. To evaluate the thermal shock behaviour, samples containing Vickers indentations were tested in quenching and the growth of the radial cracks was determined. Glycerine was chosen as the quenching fluid to minimize the effect of boiling on the heat transfer coefficient and to assure mild heat transfer conditions. A reference alumina material and a composite alumina + 10 vol% AlN with similar microstructural characteristics were prepared. The materials properties controlling the thermal shock response of the materials, i.e. Young's modulus, fracture toughness, coefficient of thermal expansion and thermal conductivity were determined and the theoretical relationships for the growth of the cracks were calculated and compared with the experimental results. The combination of thermal and mechanical properties of the composite leads to an improvement in the thermal shock resistance of alumina in terms of critical temperature increment ( $\cong 50\%$ ) which is in agreement with the value predicted by theory.

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# 1. Introduction

Monophase dense alumina materials are well known for their advantageous properties, such as wear and deformation resistance and chemical inertness, but they lack thermal shock resistance.

The thermal stresses developed in a piece of elastic material subjected to a sudden temperature change are proportional to the product of the thermal expansion coefficient,  $\alpha$ , and the Young's modulus, *E*, of the material. As a consequence, the classical parameter, *R*, that describes the thermal shock resistance of brittle materials in terms of the maximum temperature increment that the material can stand without fracture is given by<sup>1</sup>:

$$R = \frac{\sigma_{\rm f}}{E \cdot \alpha} \cdot (1 - \nu) \tag{1}$$

where  $\sigma_f$ , *E*, *v* and  $\alpha$  are the strength, the Young's modulus, the Poisson's coefficient and the coefficient of thermal expansion of the material.

Under mild thermal shock conditions, the response of the material is determined also by its thermal conductivity, k, and the classical parameter R' is given by:

$$R' = k \cdot \frac{\sigma_{\rm f}}{E \cdot \alpha} \cdot (1 - \nu) \tag{2}$$

In monophase dense ceramics it is difficult to increase this parameter because  $\alpha$  is almost fixed by phase composition and the potential ways to decrease *E*, e.g., increase porosity or microcrack density, are usually accompanied by a decrease in strength and thermal conductivity.

Two phase materials will present better thermal shock resistance than the matrix, in terms of the maximum temperature increment that the material can stand without fracture,  $\Delta T_c$ , if the second phase has lower product E· $\alpha$  and/or higher thermal conductivity than the matrix. Strength values can be maintained in the

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composite by means of a strict microstructural control, or even increased due to the toughening effect of the second phase.

In previous works,<sup>2,3</sup> we have used this approach to develop alumina-mullite and alumina-aluminium titanate composites with improved thermal shock response. The level of improvement in terms of  $\Delta T_c$  agreed qualitatively with theoretical predictions. For an alumina-10 vol% mullite material,<sup>2</sup> the increase in the critical temperature increment ( $\approx$ 12%) was due solely to the decrease of the product  $E \cdot \alpha$ , whereas for a composition alumina-10 vol% aluminium titanate,<sup>3</sup>  $\Delta T_c$  was further increased ( $\approx$ 30-40%) due to the toughening effect of the second phase. Moreover, the hardness and strength of alumina were maintained in these composites.

Using a similar approach, alumina–metal<sup>4–7</sup> and alumina–tungsten carbide<sup>8</sup> composites with improved thermal shock resistance under severe thermal shock conditions, i.e. quenching in water at room temperature were also developed. Critical temperature increments higher than those of alumina were reported for alumina–molybdenum (20 vol%),<sup>6</sup> alumina–iron (20 vol%),<sup>4,5</sup> alumina–copper (5 vol%)<sup>7</sup> and alumina–tungsten carbide<sup>8</sup> composites, and related to the reduction in Young's modulus and the increase in thermal conductivity<sup>4–7</sup> and the increase in toughness.<sup>7,8</sup>

Aluminium nitride is another possible ceramic second phase to improve the thermal shock resistance of alumina. There is not general agreement about the compatibility of AlN and alumina.<sup>9–12</sup> In particular, there is uncertainty about the formation of AlON as a function of temperature and atmosphere, due to the fact that stability in this system is highly dependent on oxygen and nitrogen partial pressures and both are determined by the gas flow rate and the furnace design, i.e. the presence of graphite. Moreover, the physical characteristics of the powders will determine the microstructural evolution during the low temperature treatments and, consequently, the accessibility of the atmospheric gas to the bulk of the compacts at higher temperature. Lastly, differences in the chemical composition of the powders, in particular the nature and amount of impurities, will also determine the phase equilibrium relationships.

Using thermodynamic data, Willems et al.<sup>12</sup> showed that AlON is stable only within a small region of oxygen and nitrogen pressures and that it is not stable below  $1640\pm10$  °C. Even though, several authors have found AlON at temperatures as low as 1600 °C,<sup>13</sup> most experiments, performed in graphite furnaces and under flowing N<sub>2</sub>, indicate that the stable phases below 1700 °C are AlN and Al<sub>2</sub>O<sub>3</sub>, while  $\gamma$ -AlON appears first at 1700 °C.<sup>9–11</sup>

Reported values of Young's modulus (350 GPa at 25 °C)<sup>14,15</sup> and coefficient of thermal expansion  $(\alpha_{25-1000 \ \circ C} = 5.6-5.7 \times 10^{-6\circ} C^{-1})^{14}$  for high purity and dense AlN materials are lower than those of dense

alumina<sup>3</sup> ( $E=405\pm2$  GPa at 25 °C, and  $\alpha_{25-600}$  °C = 9.4±0.2×10<sup>-6°</sup>C<sup>-1</sup>). For aluminium oxynitride spinel materials, stiffness is lower (E=307-330 GPa,<sup>16,17</sup> at 25 °C), and thermal expansion is slightly higher ( $\alpha_{20-1000}$  °C  $\approx$ 7.0–7.4×10<sup>-6°</sup>C<sup>-1</sup>)<sup>18</sup> than for aluminium nitride.

For pure AlN single crystal, theoretical thermal conductivity along the *c*-axis was predicted to be 319 W/ mK.<sup>19</sup> While values as high as 285 W/mK have been reported for very pure single crystals,<sup>19</sup> substantial variability ( $\approx$ 70–260 W/mK at 25 °C) exists for polycrystalline dense AlN materials depending on composition, i.e. second phases and oxygen content, and microstructure.<sup>14,20–26</sup> The thermal conductivity of aluminium oxynitride spinels is much lower ( $\approx$ 9–11 W/ mK<sup>18,27</sup> at 25 °C), and in the same range of that of the highly conductive oxides.

Even though the product E· $\alpha$  for AlN and AlON is lower than that for alumina, the reported thermal shock behaviour of hot pressed AlN and AlON ceramics under severe thermal conditions (25 °C water quench test) was similar to that of alumina,<sup>15,18</sup> which can be attributed to the relatively poor mechanical properties of the materials tested ( $K_{IC}\approx 2-3$  MPam<sup>1/2</sup>,  $\sigma_f\approx 200-350$ MPa) compared with those for similarly processed aluminas ( $K_{IC}\approx 3-4$  MPam<sup>1/2</sup>,  $\sigma_f\approx 400-500$  MPa).<sup>28,29</sup> However, the thermal shock behaviour of AlN under mild thermal conditions, where the beneficial effect of the higher thermal conductivity is expected, should be much better than that of alumina, as described by *R'* [Eq. (2)].

In this work, the possibility of improving the thermal shock behaviour of alumina under mild thermal conditions by the addition of submicron-sized and homogeneously dispersed AlN particles is studied. Processing conditions were adjusted in order to avoid the formation of AlON and, consequently, the decrease in the thermal conductivity of the composite. In principle, two different factors might play a role in the thermal response of the composite. First, from the above mentioned data, the product  $E \cdot \alpha$  for AlN ( $\approx 1.7 \text{ MPa}^{\circ}\text{C}^{-1}$ ) is about half of that for alumina ( $\approx 3.8 \text{ MPa}^{\circ}\text{C}^{-1}$ ). Second, the dispersion of a highly conductive phase would improve the thermal conductivity of the alumina matrix, as observed for vacuum sintered cordierite (92-66 vol%)-AlN (8-34 vol%) composites.<sup>30</sup> Moreover, AlN might have some toughening effect on the composite due to its thermal expansion mismatch with alumina, as observed for alumina-aluminium titanate.<sup>3</sup>

# Thermal shock test

In order to reduce the statistical effects associated to the quenching-strength method, indented samples have been successfully used by different authors to evaluate the relative thermal shock resistance of a number of dense ceramics.<sup>2,3,31–37</sup> In this work, Vickers indented samples were tested in quenching and the growth of the radial cracks as a function of the severity of the thermal shock was analyzed to evaluate the thermal shock behaviour. A modification of the method initially proposed by Osterstock et al.<sup>33,34</sup> to analyze the evolution of cracks with different sizes under the same thermal shock conditions, i.e. temperature increment,  $\Delta T$ , was used in order to evaluate the thermal shock resistance of the materials in terms of the applied temperature increment.<sup>2,3</sup> Even though this simplified analysis does not consider the existent differences between the thermal states of the surface and the bulk of the specimens,<sup>32,35</sup> it describes essentially the crack growth versus temperature increment curve and accounts for the effect of the materials properties on the thermal shock response.

The analysis of the symmetric crack pattern that appears at the corners of the Vickers indentations in brittle materials<sup>38–40</sup> shows that the stress intensity factor,  $K_0$  at the tip of a Vickers indentation flaw of length,  $C_0$ , formed due to the effect of a load, P, is given by:

$$K_0 = K_c = \chi_r \cdot P \cdot C_0^{-3/2}$$
(3)

where  $K_{\rm C}$  is the fracture toughness of the material and  $\chi_{\rm r}$  is known as the residual stress factor and is a function of the Young's modulus and the hardness of the material and the geometry of the contact between the indenter and the surface of the material.

When a specimen containing a Vickers indentation performed at a load *P* is quenched, a surface thermal stress,  $\sigma_{\rm T}$ , dependent on the temperature difference of the thermal shock,  $\Delta T$ , is produced. Therefore, the residual stress associated with the indentation is supplemented by the additional stress,  $\sigma_{\rm T}$ . For the indentation crack to advance under equilibrium conditions, the total stress intensity factor has to be equal to the fracture toughness of the material,  $K_{\rm C}$ , and the following relation between the thermal stress and the crack size is obtained:<sup>34,35</sup>

$$\sigma_{\rm T} = \frac{K_{\rm c}}{\left(\Omega \cdot \pi \cdot C\right)^{1/2}} \cdot \left(1 - \frac{\lambda \cdot \chi_{\rm r} \cdot P}{K_{\rm c} \cdot C^{3/2}}\right) \tag{4}$$

where  $(\Omega \pi C)^{1/2}$  is a geometrical factor for the stress intensity field of a half penny surface crack and  $\lambda$  $(0 \leq \lambda \leq 1)$  describes the extent of relief of residual indentation stress resulting from the heating/quenching operation.

Taking into account Eq. (3), the thermal stress can be expressed in terms of the initial crack length  $C_0$ :

$$\sigma_T = \chi_r \cdot P \cdot \left(C_0^{-3/2} - \lambda \cdot C^{-3/2}\right) / (\Omega \cdot \pi \cdot C)^{1/2}$$
(5)

The initial indentation crack of size  $C_0$  experiences stable crack growth, under the applied thermal stress up to a critical size  $C_c^{33,34}$ :

$$C_{\rm c} = \left(\frac{4 \cdot \lambda \cdot \chi_{\rm r} \cdot P}{K_{\rm c}}\right)^{2/3} \tag{6}$$

The associated stress value when  $C = C_c$  is:

$$\sigma_{\rm c} = 0.47 \cdot K_{\rm c}^{4/3} / (\pi \cdot \Omega)^{1/2} \cdot (\lambda \cdot \chi_{\rm r} P)^{1/3}$$
(7)

On this basis, Osterstock et al.<sup>33,34</sup> proposed and tested an experimental method to evaluate the relative performance of materials submitted to thermal shock. The samples were indented using different indentation loads, P, and subjected to a fixed temperature increment,  $\Delta T$ . The curves of relative increase of crack length,  $C/C_0$ , vs indentation load reflected the empirical ranking of thermal shock resistance of the materials tested.

In this work, a different procedure was used<sup>2,3</sup> in order to compare the materials in terms of the critical temperature increment. If different samples of a material are indented with the same load and subjected to increasing temperature increments,  $\Delta T$ , the pattern of crack growth is divided into three regimes: At low  $\Delta T$ , no significant crack growth can be detected; in a medium  $\Delta T$  interval the crack growth is stable; and from a larger  $\Delta T$  value, the crack grows unstably.<sup>2,3,35</sup>

For increasing temperature increments, under mild and constant heat transfer conditions and convection being the main heat transfer mechanism, the thermal stress will increase according to:<sup>1,41</sup>

$$\sigma_{\rm T} = E \cdot \alpha \cdot \Delta T \cdot \frac{f(\beta)}{(1-\nu)} \tag{8}$$

where E,  $\alpha$  and  $\nu$  are the Young's modulus, the thermal expansion coefficient and the Poisson's coefficient of the material.  $f(\beta)$  stands for geometry and heat transfer conditions and is a function of the Biot modulus,<sup>41</sup>  $\beta = ah/k$ , with h being the surface heat transfer coefficient, k the thermal conductivity of the material and a the characteristic heat transfer length.

From Eqs. (5) and (8):

$$\Delta T = \chi_{\rm r} \cdot P \cdot \left( C_0^{-3/2} - \lambda \cdot C^{-3/2} \right) / (\Omega \cdot \pi \cdot C)^{1/2}$$
$$\cdot \frac{(1-\nu)}{E \cdot \alpha} \cdot \frac{1}{f(\beta)}$$
(9)

The theoretical relationships derived from Eq. (9) for the monophase alumina and the composite have been calculated and compared to the experimental thermal shock results for the evolution of indentation cracks of equal initial crack length with increasing  $\Delta T$ .

Water at room temperature is one of the fluids most used as quenching medium in thermal shock tests. Nevertheless, this medium presents serious problems when used to test alumina-based dense materials, as they usually fail when cooled from temperatures around 200-300 °C. Samples at these temperatures originate the boiling of the water when they enter in contact with it, and, as a consequence, the surface heat transfer coefficient, *h*, experiences large variations, up to 2–3 orders of magnitude, for  $\Delta T = 200-300$  °C.<sup>42</sup>

To minimize the effect of boiling on the heat transfer coefficient and, consequently, its variability in the range of testing temperatures, glycerine was chosen as quenching fluid due to its relatively high ebullition point ( $\approx 290 \ ^\circ$ C).<sup>43</sup> Moreover, the high viscosity ( $\approx 900 \ \text{mPa} \cdot \text{s})^{43}$  and relatively low thermal conductivity ( $\approx 0.3 \ \text{W/m}^\circ\text{C})^{43}$  of this medium assured mild heat transfer conditions to account for the effect of AlN on the thermal conductivity of alumina.

#### 2. Experimental

Commercial Al<sub>2</sub>O<sub>3</sub> (Condea HPA05, Germany) and AlN (ACST, grade B, Germany) were used as starting powders, with a mean particle size (laser diffraction, Malvern Mastersizer S, UK) of 0.4  $\mu$ m for alumina and 2.1  $\mu$ m for aluminium nitride and a specific surface area (N<sub>2</sub> adsorption, BET, Monosorb, Quantachrome, USA) of 9.5 and 5.3 m<sup>2</sup>/g, for alumina and aluminium nitride, respectively. In order to decrease the AlN particle size, the powder was attrition milled with alumina balls for 10 h, to reach particle sizes smaller than 500 nm and a specific surface area of 14.8 m<sup>2</sup>/g. The mixture Al<sub>2</sub>O<sub>3</sub>–10 vol% AlN (A10AlN) was attrition milled in isopropyl alcohol with alumina balls for 1 h, dried and sieved (63 µm).

The alumina + 10 vol%AlN mixture was hot pressed at 1550 °C and 40 MPa, for 30 min in a nitrogen atmosphere. Green compacts of alumina were obtained by isostatic pressing at 200 MPa and sintering was performed in an electric furnace at 1500 °C for 1 h. The monophase alumina thermal treatment was adjusted to obtain similar grains sizes as those of the alumina grains in the composite. The obtained materials will be denoted as  $A_M$  and A–10AlN, for the reference alumina and the composite, respectively.

The crystalline phases in the sintered composite were analysed by X-ray diffraction (Siemens, D-5000, Germany) using the files ASTM10-173 and 24-1133, for alumina and aluminium nitride, respectively.

Microstructural characterisation was by scanning electron microscopy (SEM, Zeiss DSM40, Germany) on fracture surfaces and on polished and thermally etched (1450 °C-1 h in a graphite furnace in flowing Nitrogen) surfaces, as well as by transmission electron microscopy (TEM, Jeol 2000FX, Japan). For TEM studies, specimens were thinned by mechanical polishing, dimpling and ion milling; images were obtained at 200 kV.

The average coefficient of thermal expansion between room temperature and 600 °C was determined in a Linear Voltage Displacement Transducer dilatometer (Adhamel Lomargy, France) with a quartz rod. Three samples of 12 mm length and  $4 \times 5 \text{ mm}^2$  cross section were diamond cut from the sintered blocks and tested for each material. Given values are the average of the three measurements and errors relate to the maximum and minimum obtained values.

Thermal diffusivity was determined from room temperature up to 400 °C in a laser flash equipment (Holometrix Thermal Flash, USA), using thin discs (diameter = 12.7 mm, thickness = 1 mm) diamond drilled and cut from the blocks. The specific heat was evaluated by comparison with a standard of alumina. Thermal conductivity, k, was calculated from these data and the density of the materials, determined by the Archimedes' method using water. Reported values for thermal conductivity are the average of three determinations and errors relate to the maximum and minimum obtained values.

Young's modulus of the reference alumina was determined by the resonance frequency method on parallelepiped bars ( $4 \times 5 \times 50 \text{ mm}^3$ ), diamond cut from the sintered blocks, tested in flexure, using a commercial apparatus (Lemmens Grindosonic MK5 Industrial, Belgium). Given values are the average of five bars, and errors are the standard deviation of the measurements. For the composite, Young's modulus was calculated using Voight equation:

$$E_{\rm c} = v_{\rm A} \cdot E_{\rm A} + v_{\rm AIN} \cdot E_{\rm AIN} \tag{10}$$

where  $v_A$  and  $v_{AIN}$  are the volumetric fractions of alumina and AlN, respectively, and  $E_A$  and  $E_{AIN}$  are the Young's modulus of alumina and AlN, respectively. For  $E_A$ , the experimental value for the reference alumina (Table 1) and for AlN the experimental value reported for dense AlN materials (350 GPa)<sup>14,15</sup> were used.

Hardness, HV, and toughness,  $K_{\rm C}$ , were determined from Vickers indentation tests (Leco Hardness Tester, Japan) on polished surfaces (1 µm), with indentation loads of 98 N during times of 10 s. No chipping associated with any indentation was detected and the indentations that exhibited secondary radial cracking were rejected. The principal diagonals of the plastically deformed zone and the length of the radial cracks were measured using a reflected light optical microscope (RLOM, Carl-Zeiss D-7082, Germany) to calculate HV and K<sub>c</sub> values, respectively. The equation proposed by Miranzo and Moya<sup>44</sup> was used to calculate the constant  $\chi_r$  [Eq. (3)] and K<sub>c</sub>. Reported values are the average of ten determinations and errors are the standard deviations.

To fabricate small discs (diameter = 11.9, thickness = 3 mm) to be the thermal shock samples, cylinders were diamond drilled and embedded in a resin. One of the flat faces of the cylinder was diamond polished, down to 1  $\mu$ m, and subsequently a slice of 3 mm width was cut from the cylinder with a fine diamond disc to obtain the

Table 1 Density and thermo-mechanical properties of the reference alumina and the composite A10AlN (alumina + 10 vol% AlN)

	Alumina	A-AlN
Density (g/cm <sup>3</sup> )	$3.94 \pm 0.02$	$3.90 \pm 0.01$
% of theoretical	98.6	99.5
$\alpha_{200-600 \ ^{\circ}C} (\times 10^{-6} K^{-1})$	$9.4 \pm 0.2$	$9.1 \pm 0.2$
Young's modulus (GPa)	$405 \pm 2$	373
Vicker's Hardness (GPa)	$19.4 \pm 0.3$	$17.9 \pm 0.2$
Toughness, Kc (MPam <sup>1/2</sup> )	$3.8 \pm 0.2$	$4.5 \pm 0.1$

test sample. During the cutting process parallelism of the surfaces was maintained and dimensions were carefully checked. Previous to the thermal shock test, a Vickers indentation (98 N for 10 s) was introduced in the centre of the polished face of the discs. The indentations formed cracks of similar sizes in both materials  $(C_0 = 203 \pm 10 \text{ and } 183 \pm 8 \ \mu\text{m}$  for  $A_M$  and A–10AlN, respectively). For the thermal shock tests, each sample was fixed in an aluminium sample holder with open bottom. To avoid stress concentration at the edges and at the contact between the holder and the indented area of the specimen, a fibre blanket was placed around the lateral surface of the disc and extended 1 mm into the indented surface.

The holder was heated in a vertical electrical furnace using 5 °C/min as the heating rate and 20 min as the stabilisation time before quenching in a glycerine bath at 40 °C. The indented faces of the discs were placed facing the bottom of the holder and thus were the first to contact the cooling medium. The size of the indentation cracks was measured before and after the thermal shock test using a reflected light optical microscope (Carl-Zeiss D-7082, Germany).

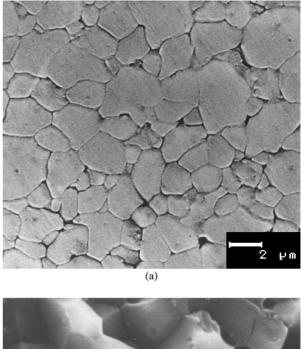
# 3. Results and discussion

#### 3.1. Microstructure and physical properties

In the composite material, only  $Al_2O_3$  and AlN were identified by XRD, and no evidence of AlON was detected, as expected from the sintering schedule used ( $T < 1640 \ ^{\circ}C$ ), in agreement with the thermodynamic data supplied by Willems et al.<sup>12</sup> and the phase equilibrium diagram proposed by McCauley and Corbin<sup>10,11</sup>

Table 1 summarises the densities and the thermomechanical properties of both materials and Figs. 1 and 2 show the microstructure of the composite. Aluminium nitride could not be differentiated from alumina by SEM of the polished surfaces (Fig. 1a); in these micrographs equi-dimensional grains of sizes between 0.5 and 3  $\mu$ m were observed (Fig. 1a). Conversely, in the fracture surfaces (Fig. 1b) small (< 500 nm) and whiter grains were found located at triple points and inside the alumina grains. These small grains were identified by electron diffraction in the selected area diffraction mode (TEM) as aluminium nitride (Fig. 2). Both diffraction patterns accompanying images in Fig. 2 correspond to the (0001) plane of the AlN hexagonal structure (a=0.3111 nm, c=0.4979 nm, JCPDS 25–1133). The intra-granular AlN particles (<200 nm, Fig. 2a) were smaller than those located at triple points and at grain boundaries (200–500 nm, Figs. 1 and 2b). Only alumina and aluminium nitride were observed by TEM, in agreement with the XRD results.

The contribution of AlN to the thermal expansion of the composite can be evaluated using the Turner's expression:<sup>45</sup>



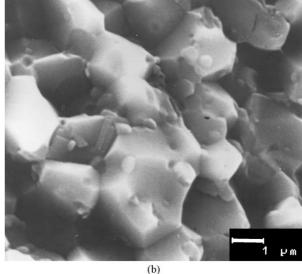
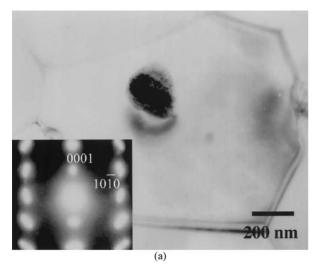


Fig. 1. Microstructure of the A10AlN (Alumina+10 vol%AlN) composite. (a) SEM of a polished and thermally etched surface. Equidimensional grains of sizes between 0.5 and 3  $\mu$ m are observed. (b) SEM of a fracture surface. Small (<500 nm) and whiter grains are found located at triple points, grain boundaries and inside the alumina grains.



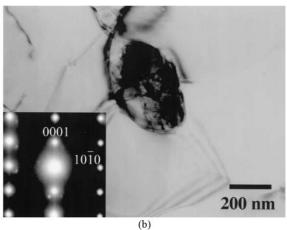


Fig. 2. TEM micrographs of the A10AlN composite. The two locations of the AlN particles are observed, as shown by the diffraction diagrams. (a) AlN particle embedded in an alumina grain. (b) AlN particle embedded located at a triple point.

$$\alpha_{\rm c} = \frac{\frac{\alpha_{\rm A} \cdot K_{\rm A} \cdot \nu_{\rm A}}{\rho_{\rm A}}}{\frac{K_{\rm A} \cdot \nu_{\rm A}}{\rho_{\rm A}}} + \frac{\frac{\alpha_{\rm AIN} \cdot K_{\rm AIN} \cdot \nu_{\rm AIN}}{\rho_{\rm AIN}}}{\frac{K_{\rm AIN} \cdot \nu_{\rm AIN}}{\rho_{\rm AIN}}}$$
(11)

where  $\alpha_c$  is the thermal expansion coefficient of the composite (Table 1), and  $\alpha$ ,  $\rho$ , v and K are the thermal expansion coefficient, the density, the volume fraction and the bulk modulus of alumina (A) and aluminum nitride (AlN), respectively.

The bulk modulus of the material, *K*, is given by:

$$K = \frac{E}{3 \cdot (1 - 2 \cdot \upsilon)} \tag{12}$$

where E is the Young's modulus and v is the Poisson coefficient of the material.

Introducing the values of the properties summarised in Table 1, and assuming coincident values of the Poisson coefficient for both materials, a value of  $5.8 \times 10^{-6\circ}$ C<sup>-1</sup> is obtained for the thermal expansion coefficient of the AlN particles dispersed in the alumina matrix. This value is coincident with that reported for high purity and dense aluminium nitride ceramics.<sup>14</sup>

Hardness of the composite (Table 1) is slightly lower than that of alumina, as expected due to the addition of second phase with lower hardness (H $\approx$ 14–15 GPa for dense AlN<sup>17,18</sup>), but still in the range of the hardness of dense aluminas. This value is coincident with that reported for hot pressed alumina–AlN composites with similar composition ( $\approx$ 18GPa).<sup>14</sup>

Toughness of the composite is higher than that of alumina, as reported for an alumina (90 vol%) + AlON (10 vol%) composite ( $K_{IC}$ =4.35 MPam<sup>1/2</sup>, SENB).<sup>29</sup> Conversely, larger amounts (>20 vol%) of AlON were reported to have no effect<sup>29</sup> or even decrease<sup>46</sup> the toughness of the alumina matrix. The exact mechanism acting during fracture of the alumina–AlN composite studied here has to be further analysed but, in principle, microcracking and/or crack bridging might occur due to the compressive stresses developed around the second phase particles originated by the thermal expansion mismatch between alumina and AlN. The presence of radial microcracks around some AlN particles in the samples observed by TEM (Fig. 3) is evidence of these stresses.

Thermal conductivity as a function of temperature is plotted in Fig. 4. Values are higher for the composite across the whole temperature range tested. Room temperature value is significantly higher than the value ( $\approx 18$ W/mK)<sup>14</sup> reported for alumina–AlN composites of similar composition hot pressed at 1700–1800 °C. This difference can be attributed to the presence of significant amounts of AlON in the latter materials due to the higher sintering temperatures.

The contribution of the AlN particles to the thermal conductivity of the composite,  $k_c$ , can be calculated by fitting the Maxwell relationship for composite materials:<sup>45</sup>

$$k_{\rm c} = k_{\rm A} \frac{1 + 2 \cdot v_{\rm AIN} \cdot (1 - k_{\rm A}/k_{\rm AIN})/(2k_{\rm A}/k_{\rm AIN} + 1)}{1 - v_{\rm AIN}(1 - k_{\rm A}/k_{\rm AIN})/(k_{\rm A}/k_{\rm AIN} + 1)}$$
(13)

Where k is the thermal conductivity, and v, A, and AlN stand for the volumetric fraction, alumina and aluminium nitride, as previously.

In order to eliminate the effect of porosity differences (Table 1) the thermal conductivity values for dense materials were calculated from the experimental values (Fig. 4) using the Klemens approximation:<sup>47</sup>

$$\frac{kp(\text{porous})}{kp(\text{dense})} = 1 - \frac{4}{3}\phi \tag{14}$$

Were,  $\phi$  is the pore volume fraction.

Introducing room temperature values, a thermal conductivity value of about 50 W/mK is obtained for the

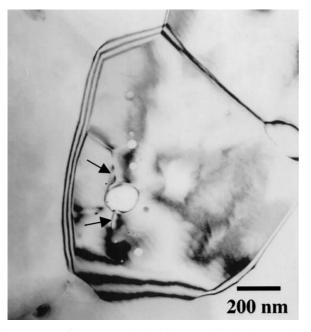


Fig. 3. TEM of the composite A10AlN material showing an intragranular AlN particle with associated radial microcracks signalled by arrows.

aluminium nitride particles, which is in the range of those (69.5 W/mK<sup>22</sup> and 62 W/mK<sup>23</sup> for AlN-Y<sub>2</sub>O<sub>3</sub>-Al<sub>2</sub>O<sub>3</sub> sintered at 1850 and 1750 °C, respectively) reported for dense AlN materials with high oxygen contents in the AlN grains ( $\geq 1.5$  wt.%).

The evolution of thermal conductivity of the composite with temperature, plotted in Fig. 4, as well as the values of the coefficient of thermal expansion and hardness demonstrate that no or few microcracks are present in this material. Therefore, it is justified the use of the Voight expression [Eq. (10)] to calculate the Young's modulus of this material. The microcracks observed by

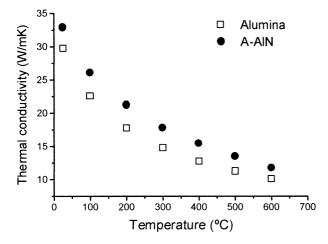


Fig. 4. Thermal conductivity as a function of temperature for the reference alumina material and the composite A10AlN (alumina +10 vol% of AlN). The size of the symbols accounts for experimental variability.

TEM might be formed during the preparation process of the samples due to the high residual stresses present in the material.

## 3.2. Thermal shock

In Fig. 5 the experimental results for thermal shock are plotted. There is a threshold temperature increment for crack propagation,  $\Delta T_0$ , for each material,  $\Delta T_0 \approx 180$  °C for alumina and 300 °C <  $\Delta T_0$  < 400 °C for the composite. Collin and Rowcliffe observed this temperature threshold for a number of different materials quenched in water at room temperature, and concluded that it was due to the large differences between the heat transfer coefficients of surfaces at temperatures lower and higher than 100 °C.<sup>35</sup> This argument can not be applied to the experiments reported here in which the heat transfer coefficient is low and essentially constant across the whole range of temperature increments considered.

Taking into account Eq. (4), the existence of the threshold temperature might be due to stress relief  $(\lambda < 1)$ . As no chipping was observed in any of the tested samples, no stress relief occurred by additional lateral cracking. Moreover, no stress relief by mass transport is expected during heating of alumina based materials due to the relatively low temperatures involved. Lee et al.<sup>36</sup> attributed the existence of this temperature threshold in indentation-quench tests of silicon nitride to post-indentation relaxation of residual contact stresses before quenching. On the other hand, the analysis by Fett<sup>48</sup> indicates that the action of a crack opening force, such as the thermal stress considered here, will also lead to the relaxation of the residual stresses.

Stable crack growth occurred in all samples tested using  $\Delta T$  lower or equal to 300 and 500 °C for alumina

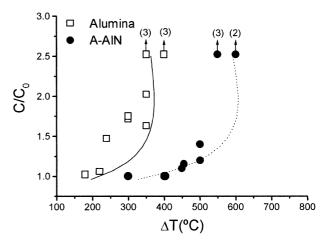


Fig. 5. Thermal shock test results for the reference alumina material and the composite A10AlN (alumina + 10 vol% of AlN). Relative crack growth ( $C/C_0$ ) versus temperature increment. Lines corresponds to the theoretical relationships [Eq. (9)] calculated from the properties of the materials.

and A-10AlN, respectively ( $C/C_0 \le 1.7$ , Fig. 5). The radial cracks traversed more than 50% of the samples tested for  $\Delta T$  larger or equal to 350 and 550 °C for alumina and the composite, respectively. Therefore, the critical temperature increments,  $\Delta T_C$ , are located between 300 and 350 °C for alumina and between 500 and 550 °C for A-10AlN, giving the following ratio between the critical temperature increments for the reference alumina,  $\Delta T_C$  A and the composite,  $\Delta T_C$  A–10AlN:

$$1.4 \leqslant \frac{\Delta T_{\rm CA-10AIN}}{\Delta T_{\rm CA}} < 1.8 \tag{15}$$

This level of improvement of the thermal shock of alumina in terms of  $\Delta T_c$  is comparable to that found for alumina–iron (20 vol%)<sup>5,6</sup> and alumina–copper (5 vol%)<sup>7</sup> composites obtained by hot pressing for which the ratio was 1.5, and higher than those found for alumina–mullite<sup>2</sup> and alumina–aluminum titanate composites<sup>3</sup> tested under the same thermal conditions as those used here.

The value of  $\lambda$  can be estimated from the experimental results from Fig. 5. Taking into account Eqs. (3) and (6), the following relationship between the initial and the critical crack sizes and is obtained:

$$C_{\rm c} = (4 \cdot \lambda)^{2/3} \cdot C_0 \tag{16}$$

From Eqs. (9) and (16) and making  $\Delta T_0 = \Delta T$ , for  $C = C_0$  and  $\Delta T_C = \Delta T$ , for  $C = C_C$ :

$$\frac{\Delta T_0}{\Delta T_C} = 2.13 \cdot (1 - \lambda) \cdot \lambda^{-1/3} \tag{17}$$

Introducing the experimental values of  $\Delta T_0$  and  $\Delta T_C$  for the reference alumina and the composite,  $\lambda \approx 0.7$  is obtained for both materials, which indicates that a stress release of about 30% occurs. This value is in the range of those calculated from experimental data from other authors  $(0.6-0.7)^{49,50}$  and for alumina–mullite and alumina–aluminium titanate composites tested under similar thermal conditions ( $\approx 0.7$ ).<sup>3,4</sup> For this  $\lambda$  value,  $C_C/C_0$  in Eq. (16) is about 2, in agreement with the fact that, relative crack lengths  $C/C_0 \leq 1.7$  were determined in all samples in which stable fracture was observed.

The experimental values plotted in Fig. 5 can be compared with the theoretical relationships derived from Eq. (9). For Biot modulus lower than 5, the function  $f(\beta)$  in Eqs. (8 and 9) can be approximated by:<sup>41</sup>

$$f(\beta) = \frac{\beta}{B} \tag{18}$$

where B is a constant between 3 and 4.

The Biot modulus for the experiments considered here, evaluated from the value of the heat transfer coefficient determined by Sato et al.<sup>51</sup> for quenching of tetragonal zirconia in glycerine ( $\approx 1000 \text{ W/m}^2\text{K}$ ), the size of the samples (3 mm) and the thermal conductivity values from Fig. 4, is  $\beta < 1$ , therefore, Eq. (18) can be applied.

In Fig. 5, the theoretical relationships derived from Eq. (9) and data from Table 1 and Fig. 4 are plotted. B=3.25, as proposed by Manson,<sup>41</sup>  $h=1000 \text{ W/m}^2\text{K}$ ,<sup>51</sup> and the thermal conductivity of the bulk of the samples ( $T\approx300 \text{ °C}$ ) have been used to evaluate the Biot modulus for each material. From this plot,  $\Delta T_{\text{C}}=370$ , 585 °C are derived for alumina and the composite, respectively, which ratio (1.58) is very close to the experimental one [Eq. (15)].

#### 4. Conclusions

Small concentrations (10 vol%) of aluminium nitride added to alumina in the form of submicron sized and homogeneously distributed particles lead to a decrease in Young's modulus and thermal expansion coefficient of the alumina matrix and an increment in toughness. Thermal conductivity is increased when the presence of AlON phases is avoided by low temperature thermal treatment in a graphite furnace in flowing nitrogen. The combination of thermal and mechanical properties thermal conductivity, thermal expansion coefficient, Young's modulus and toughness—of the composite leads to an improvement of the thermal shock resistance of alumina in terms of the critical temperature increment ( $\cong$  50%) which agrees with that derived from theoretical calculations.

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